

1/16/09

Plane wave solutions

①

recall: $\nabla^2 \vec{E} - \mu_0 \epsilon \frac{\partial^2 \vec{E}}{\partial t^2} = 0$

plane wave - no boundary conditions \Rightarrow infinite waves (unphysical)

use separation of variables

$$\vec{E} = \vec{E}_0 E(\vec{r}) E(t)$$

\vec{E}_0 has ampl. and polarization $|\vec{E}_0| = E_0$

$$\vec{E}_0 E(t) \nabla^2 E(\vec{r}) - \mu_0 \epsilon \vec{E}_0 E(\vec{r}) \frac{\partial^2 E(t)}{\partial t^2} = 0$$

$$\frac{\nabla^2 E(\vec{r})}{E(\vec{r})} - \mu_0 \epsilon \frac{\partial^2 E(t) / \partial t^2}{E(t)} = 0$$

for this eqn. to be valid, each term must be a constant for all space and time

$$\frac{\nabla^2 E(\vec{r})}{E(\vec{r})} = -k^2 \Rightarrow \nabla^2 E(\vec{r}) = -k^2 E(\vec{r})$$

$$\Rightarrow E(\vec{r}) = e^{-j\vec{k} \cdot \vec{r}} \quad (+c.c.) \quad k = |\vec{k}|$$

now substitute into wave equation

$$-k^2 - \mu_0 \epsilon \frac{\partial^2 E(t) / \partial t^2}{E(t)} = 0$$

$$\Rightarrow \frac{\partial^2 E(t)}{\partial t^2} = -\frac{k^2}{\mu_0 \epsilon} E(t) = -\omega^2 E(t)$$

$$\Rightarrow E(t) = e^{j\omega t} \quad (+c.c.)$$

$$\vec{E} = \vec{E}_0 e^{j(\omega t - \vec{k} \cdot \vec{r})}$$

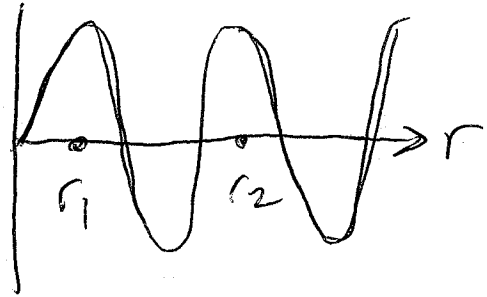
$\omega =$ radial freq. rad/s

define $\omega^2 = \frac{k^2}{\mu_0 \epsilon}$

so that $k = \omega \sqrt{\mu_0 \epsilon}$
wave number rad/m

$\vec{k} =$ wave vector dir. of phase propagation

$E(r)E(t) \quad t=0$



in space, distance between two adjacent peaks is wavelength λ (2)

$$e^{-jk r_1} = e^{-jk r_2}$$

$$= e^{-jk(r_1 + \lambda)}$$

$$\Rightarrow e^{-jk\lambda} = 1 \quad \text{or} \quad \lambda = 2\pi/k$$

$$(k = 2\pi/\lambda)$$

can do same for time to get period $T = 2\pi/\omega = 1/\nu$

Transverse waves

assume propagation along z and polarization along x

$$\vec{E}(r,t) = \hat{x} E_0 \cos(\omega t - kz)$$

$$= \frac{\hat{x}}{2} E_0 e^{-j(kz - \omega t)} + \text{c.c.} \quad \text{or take real part}$$

Find magnetic field

$$\text{use } \nabla \times \vec{E} = -\mu_0 \frac{d\vec{H}}{dt}$$

$$\nabla \times \vec{E} = \det \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix}$$

$$= -\hat{y} \left(-\frac{\partial E_x}{\partial z} \right) + \hat{z} \left(-\frac{\partial E_x}{\partial y} \right)$$

only y-component of $\vec{H} = H_y$

$$-jk E_0 e^{-j(kz - \omega t)}$$

$$= -\mu_0 \frac{dH_y}{dt}$$

$$H_y \propto e^{-j(kz - \omega t)} + \text{c.c.}$$

$$\frac{jk E_0}{\mu_0} e^{-j(kz - \omega t)}$$

$$\frac{1}{j\omega} + \text{c.c.} = H_y$$

$$H_y = \frac{k}{\mu_0 \omega} E_0 e^{-j(kz - \omega t)} + \text{c.c.}$$

$\mu_0 = 377 \Omega$
impedance

$$\frac{k}{\mu_0 \omega} = \frac{\sqrt{\mu_0 \epsilon'}}{\mu_0} = \frac{1}{\mu_0} \sqrt{\frac{\epsilon'}{\mu_0}} = \frac{1}{Z_0}$$

Poynting vector - power flow through surface (3)

$$\vec{S} = \vec{E} \times \vec{H} \quad (\text{W/m}^2)$$

\hat{S} = direction of power flow
(not necessarily = \hat{k})

using complex notation

$$\langle \vec{S} \rangle = \frac{1}{2} \text{Re} [\vec{E}_0 \times \vec{H}_0^*]$$

time average, valid w/o c.c.

~~Re[E₀]~~

$$\langle S_z \rangle = \frac{1}{2} \text{Re} [\vec{E}_0 \times \vec{H}_0^* \cdot \hat{z}]$$

$$\vec{E}_0 \times \vec{H}_0^* = \det \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ E_0 & 0 & 0 \\ 0 & \frac{E_0^*}{\eta} & 0 \end{vmatrix} = \hat{z} \frac{|E_0|^2}{\eta}$$

$$\langle S_z \rangle = \frac{1}{2} \text{Re} [\sqrt{\cdot} \cdot \hat{z}] = \frac{|E_0|^2}{2\eta} \text{ Intensity}$$

phase velocity

$$e^{-j(kz - \omega t)} = \text{const.}$$

$$\Rightarrow kz - \omega t = \text{const}$$

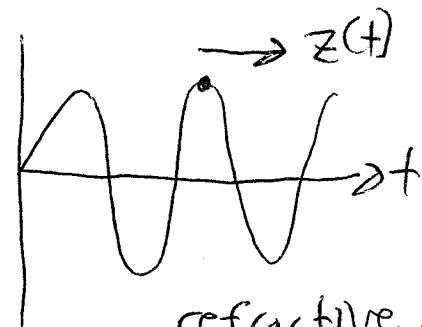
$$z(t) = \text{const} + \omega t / k$$

$$v_p = \frac{dz(t)}{dt} = \frac{\omega}{k} = \frac{1}{\sqrt{\mu_0 \epsilon}}$$

refractive index

$$n = \frac{c}{v_p} = \sqrt{\frac{\epsilon}{\epsilon_0}}$$

$$k = n k_0 \quad k_0 = \frac{2\pi}{\lambda_0} = \text{free space}$$



group velocity

(4)

velocity of envelope of superposition of waves

$$\omega_1 = \omega + \Delta\omega$$

$$\omega_2 = \omega - \Delta\omega$$

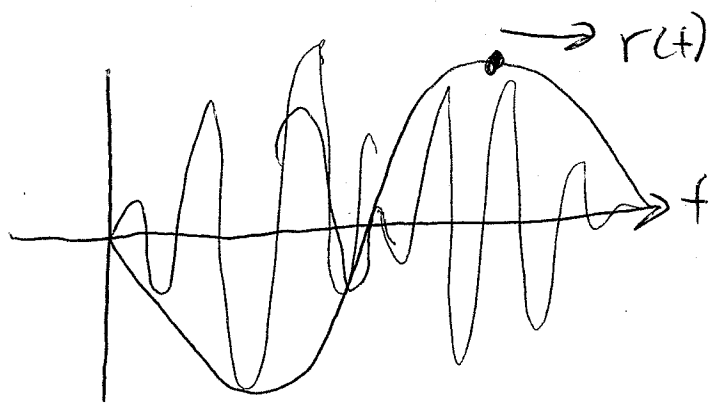
$$k_1 = k + \Delta k$$

$$k_2 = k - \Delta k$$

$$E_1 + E_2 = E_0 \left\{ \cos[(\omega + \Delta\omega)t - (k + \Delta k)z] + \cos[(\omega - \Delta\omega)t - (k - \Delta k)z] \right\}$$

$$\cos(x+y) + \cos(x-y) = 2\cos(x)\cos(y)$$

$$E_1 + E_2 = 2E_0 \cos(\omega t - kz) \cos(\Delta\omega t - \Delta k z)$$



now, envelope velocity

$$z(t) = \frac{\Delta\omega t}{\Delta k} + \text{const.}$$

$$v_g = \frac{dz(t)}{dt} = \frac{\Delta\omega}{\Delta k} \Rightarrow \frac{d\omega}{dk}$$

in free space $v_g = \frac{d\omega}{dk} = \frac{d(k \cdot c)}{dk} = c$

in dispersive medium, i.e. $n(\omega)$ not const.

$$v_g = \frac{d}{dk} \left(\frac{kc}{n} \right) = \frac{c}{n} - \frac{kc}{n^2} \frac{dn}{dk}$$
$$= \frac{c}{n} - \lambda \frac{dn}{d\lambda}$$