

ECE 3300 Maxwell's Equations in Free Space and Plane Waves

Maxwell's Equations (General – Time Domain and Phasor Forms)

(To convert from time domain to frequency/phasor domain, assume that E,H are functions of $e^{j\omega t}$. Then $d/dt = j\omega$)

$$\nabla \cdot \bar{D} = \rho_v$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} = -j\omega\mu\bar{H}$$

$$\nabla \cdot \bar{B} = 0$$

$$\nabla \times \bar{H} = \bar{J}_c + \frac{\partial \bar{D}}{\partial t} = \sigma\bar{E} + j\omega\epsilon\bar{E} = j\omega\epsilon_c\bar{E}$$

$$\epsilon_c = \epsilon_o\epsilon_r - j\sigma/\omega = \epsilon' - j\epsilon''$$

Here, J_c = conduction current density (A/m²), ϵ_c is complex dielectric constant.

Wave Equations in Charge-Free / Current-free Medium:

If the material is lossless ($\sigma=0$), then there is no charge, and no current. This is called a “perfect dielectric”. An insulator is a good or near-perfect dielectric.

$$\nabla \cdot \bar{E} = 0$$

$$\nabla \times \bar{E} = -j\omega\mu\bar{H}$$

$$\nabla \cdot \bar{H} = 0$$

$$\nabla \times \bar{H} = j\omega\epsilon_c\bar{E}$$

Your book derives the general wave equation in the time domain. So we will derive the wave equation in the frequency (phasor) domain instead

Derivation of WaveEquation:

Take $\nabla \times$:

$$\nabla \times (\nabla \times \bar{E}) = -j\omega\mu(\nabla \times \bar{H})$$

Substitute:

$$\nabla \times (\nabla \times \bar{E}) = -j\omega\mu(j\omega\epsilon_c \bar{E}) = \omega^2 \mu\epsilon_c \bar{E}$$

Vector _ property:

$$\nabla \times (\nabla \times \bar{E}) = \nabla(\nabla \cdot \bar{E}) - \nabla^2 \bar{E}$$

Substitute:

$$\nabla \times (\nabla \times \bar{E}) = \nabla(0) - \nabla^2 \bar{E} = -\nabla^2 \bar{E}$$

Combine:

$$\omega^2 \mu\epsilon_c \bar{E} - \nabla^2 \bar{E} = 0$$

$$\nabla^2 \bar{E} - \gamma^2 \bar{E} = 0$$

$$\gamma^2 = \omega^2 \mu\epsilon_c = \alpha + j\beta$$

$$\text{Lossless: } k^2 = \omega^2 \mu\epsilon = \beta^2$$

$$\nabla^2 \bar{E} - k^2 \bar{E} = 0$$

Similarly:

$$\nabla^2 \bar{H} - \gamma^2 \bar{H} = 0$$

These wave equations can be used to describe any type of wave (planar, spherical, etc.). No approximations or limitations have been made (yet).

Plane Waves are a special type of wave:

- Planar wavefront of infinite extent
- TEM: E,H are perpendicular to each other and to the direction of propagation
- Plane waves are like transmission line waves (which are also TEM). Look for the similarities.

Here are the transmission line wave equations:

$$-d^2 I(z)/dz^2 - \gamma^2 I(z) = 0$$

$$-\frac{d^2 V(z)}{dz^2} - \gamma^2 V(z) = 0$$

How do we solve problems with the wave equation:

- 1) Write wave equation (for instance in rectangular coordinates).
- 2) Separate the vector components into x,y,z. "Equate vector components", which means to do a separate equation for each vector component.
- 3) For plane wave front (flat wavefront) moving in the z-direction, d/dx and $d/dy = 0$. Substitute this into (2) above.
- 4) Verify that E_z or $H_z = 0$. PLANE WAVE HAS NO FIELD COMPONENTS in direction of propagation. It is a TEM wave like the ones we have studied in transmission lines.

General solution for a TEM wave: (such as Plane wave)

This applies to vector fields AND to each component separately.

$$\bar{E}(z) = \bar{E}^+(z) + \bar{E}^-(z) = \bar{E}_0^+ e^{-j\gamma z} + \bar{E}_0^-(z) e^{j\gamma z}$$

Now, for a plane-wave traveling in the z-direction, $E_z=0$, but E_x and E_y may not be zero. Solve for \mathbf{H} from Maxwell's equations (see text), and we obtain:

$$H_{y0}^+ = \frac{\gamma}{\omega\mu} E_{x0}^+ = \frac{1}{\eta} E_{x0}^+$$

Where η is **Intrinsic Impedance** of medium.

$\eta=377$ for air

η =real for lossless materials (E and H are perpendicular and in-phase)

η =complex for lossy materials. (E and H are perpendicular and out-of-phase)

Duality to Transmission Lines:

$E \rightarrow V$ (recall also $V = - \int E \, dl$)

$H \rightarrow I$ (recall also $I = \text{closed integral } H \, dl$)

$\eta \rightarrow Z_0$

Other Similarities:

- 1) Waves can be broken into + and – traveling waves
- 2) Propagate in z-direction as e^{-jkz} ($k=\beta$), including effect of phase change and attenuation
- 3) Waves will reflect, depending on the material (discontinuity) they hit.
 - (a) Reflection coefficient, transmission coefficient
 - (b) Standing Waves
 - (c) Smith Chart solution

Differences:

- 1) V and I are not vector quantities, E and H are. This means we can have “polarization”.