Ex:

a) Using superposition, derive an expression for $v_{\mathrm{a}, \mathrm{b}}$ that contains no circuit quantities other than $i_{\mathrm{s}}, v_{\mathrm{s}}, R_{1}, R_{2}$, and $\alpha$. Current $i_{\mathrm{x}}$ must not appear in your solution. Note: $\alpha \geq 0$.
b) Make a consistency check on your expression for $v_{\mathrm{a}, \mathrm{b}}$ by setting resistors and sources to numerical values for which the value of $v_{\mathrm{a}, \mathrm{b}}$ is obvious. State the values of resistors and sources for your consistency check, and show that your expression for $v_{\mathrm{a}, \mathrm{b}}$ is satisfied for these values. (In other words, plug the values into your expression from part (a) and show that it agrees with the value from your consistency check.)
c) Find the Thevenin equivalent circuit at terminals a and b. Express the Thevenin voltage, $v_{\mathrm{Th}}$, and Thevenin resistance, $R_{\mathrm{Th}}$ in terms of no circuit quantities other than $i_{\mathrm{s}}, v_{\mathrm{s}}, R_{1}, R_{2}$, and $\alpha$. $i_{\mathrm{x}}$ must not appear in your solution. Note: $\alpha \geq 0$.
d) Find an expression for the value of $R_{\mathrm{L}}$ connected from $\mathbf{a}$ to $\mathbf{b}$ that would absorb maximum power. Your answer must be written in terms of no circuit quantities other than $i_{\mathrm{s}}, v_{\mathrm{s}}, R_{1}, R_{2}$, and $\alpha$. Note: $\alpha \geq 0$.

SoL'n: a) In superposition, we turn on one independent source at a time, and we keep dependent sources on all the time.

Case I: Turn on $v_{\mathrm{s}}$ and turn off $i_{\mathrm{S}}$ (which becomes an open circuit).
Using the node-voltage method, we label the entire top wire to the right of $v_{\mathrm{s}}$ as $v_{\mathrm{a}, \mathrm{b}}$ and we label the bottom wire as reference. We then write an expression for $i_{\mathrm{x}}$ in terms of node-voltage $v_{\mathrm{a}, \mathrm{b}}$ and write an equation for the sum of currents out of the $v_{\mathrm{a}, \mathrm{b}}$ node.

$$
i_{\mathrm{x}}=\frac{v_{\mathrm{a}, \mathrm{~b} 1}-v_{\mathrm{s}}}{R_{1}}
$$

Now the sum-of-currents equation:

$$
\frac{v_{\mathrm{a}, \mathrm{~b} 1}-v_{\mathrm{s}}}{R_{1}}+\frac{v_{\mathrm{a}, \mathrm{~b} 1}}{R_{2}}+\alpha \frac{v_{\mathrm{a}, \mathrm{~b} 1}-v_{\mathrm{s}}}{R_{1}}=0 \mathrm{~A}
$$

We factor out $v_{\mathrm{a}, \mathrm{b}}$ and move constants to the other side of the equation:

$$
v_{\mathrm{a}, \mathrm{~b} 1}\left(\frac{1}{R_{1}}+\frac{1}{R_{2}}+\alpha \frac{1}{R_{1}}\right)=v_{\mathrm{s}}\left(\frac{1}{R_{1}}+\alpha \frac{1}{R_{1}}\right)
$$

or

$$
v_{\mathrm{a}, \mathrm{~b} 1}=v_{\mathrm{s}} \frac{\frac{1+\alpha}{R_{1}}}{\frac{1+\alpha}{R_{1}}+\frac{1}{R_{2}}}=v_{\mathrm{s}} \frac{(1+\alpha) R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

An alternative approach is to model the dependent source as a resistor, $R_{\text {eq. }}$. We observe that $v_{\mathrm{a}, \mathrm{b}}$ may then be computed by using a voltage divider:

$$
v_{\mathrm{a}, \mathrm{~b} 1}=v_{\mathrm{s}} \frac{R_{2} \| R_{\mathrm{Eq}}}{R_{1}+R_{2} \| R_{\mathrm{Eq}}}
$$

To find the equivalent resistance of the dependent source, we observe that current $i_{\mathrm{x}}$ flows through $R_{2}$ in parallel with $R_{\text {eq }}$, so we have a current divider with currents $(1+\alpha) i_{\mathrm{x}}$ in $R_{2}$ and $-\alpha i_{\mathrm{x}}$ in $R_{\mathrm{eq}}$. The ratio of the currents is the inverse of the ratio of the resistances:

$$
\frac{(1+\alpha) i_{\mathrm{x}}}{-\alpha i_{\mathrm{x}}}=\frac{R_{\mathrm{eq}}}{R_{2}}
$$

or

$$
R_{\mathrm{eq}}=-R_{2} \frac{1+\alpha}{\alpha}
$$

Now we can compute the parallel resistance of $R_{2}$ and $R_{\text {eq }}$.

$$
R_{2}\left\|R_{\mathrm{eq}}=R_{2}\right\|-\frac{R_{2}(1+\alpha)}{\alpha}=R_{2} \cdot 1 \|-\frac{1+\alpha}{\alpha}=R_{2} \frac{-\frac{1+\alpha}{\alpha}}{1-\frac{1+\alpha}{\alpha}}
$$

or

$$
R_{2} \| R_{\mathrm{eq}}=R_{2} \frac{-(1+\alpha)}{\alpha-(1+\alpha)}=R_{2}(1+\alpha)
$$

We use this result in the voltage divider, obtaining the same result as before:

$$
v_{\mathrm{a}, \mathrm{~b} 1}=v_{\mathrm{s}} \frac{R_{2} \| R_{\mathrm{Eq}}}{R_{1}+R_{2} \| R_{\mathrm{Eq}}}=v_{\mathrm{s}} \frac{R_{2}(1+\alpha)}{R_{1}+R_{2}(1+\alpha)}
$$

Case II: Turn off $v_{\mathrm{s}}$ (which becomes a wire) and turn on $i_{\mathrm{s}}$.
Using node voltage, we proceed as in Case I but have a simpler equations because $v_{\mathrm{S}}$ is off.

$$
i_{\mathrm{x} 2}=\frac{v_{\mathrm{a}, \mathrm{~b} 2}}{R_{1}}
$$

Now the sum-of-currents equation:

$$
\frac{v_{\mathrm{a}, \mathrm{~b} 2}}{R_{1}}+\frac{v_{\mathrm{a}, \mathrm{~b} 2}}{R_{2}}+\alpha \frac{v_{\mathrm{a}, \mathrm{~b} 2}}{R_{1}}-i_{\mathrm{s}}=0 \mathrm{~A}
$$

We factor out $v_{\mathrm{a}, \mathrm{b}}$ and move constants to the other side of the equation:

$$
v_{\mathrm{a}, \mathrm{~b} 2}\left(\frac{1}{R_{1}}+\frac{1}{R_{2}}+\alpha \frac{1}{R_{1}}\right)=i_{\mathrm{s}}
$$

or

$$
v_{\mathrm{a}, \mathrm{~b} 2}\left[R_{2}(1+\alpha)+R_{1}\right]=i_{\mathrm{s}} R_{1} R_{2}
$$

or

$$
v_{\mathrm{a}, \mathrm{~b} 2}=i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

An alternative approach is to model the dependent source as a resistor. We have the dependent source in parallel with $R_{1}$, and we may calculate the voltage across $R_{1}$ as $i_{\mathrm{x}} R_{1}$. This voltage is also across the dependent source, allowing us to define an equivalent resistance for the dependent source using Ohm's law:

$$
R_{\mathrm{Eq} 2}=\frac{i_{\mathrm{x}} R_{1}}{\alpha i_{\mathrm{x}}}=\frac{R_{1}}{\alpha}
$$

It is interesting to note that this equivalent resistance is different than the equivalent resistance from Case I.

Now we have current source $i_{\mathrm{s}}$ in parallel with three resistors that we can combine into one resistance, and Ohm's law gives $v_{\mathrm{a}, \mathrm{b} 2}$ in terms of current times resistance:

$$
v_{\mathrm{a}, \mathrm{~b} 2}=i_{\mathrm{s}} \cdot R_{1}\left\|R_{2}\right\| R_{\mathrm{Eq} 2}=i_{\mathrm{s}} \cdot R_{1}\left\|R_{2}\right\| \frac{R_{1}}{\alpha}
$$

or

$$
v_{\mathrm{a}, \mathrm{~b} 2}=i_{\mathrm{s}} \frac{1}{\frac{1}{R_{1}}+\frac{1}{R_{2}}+\frac{\alpha}{R_{1}}}=i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

Now we sum the two $v_{\mathrm{a}, \mathrm{b}}$ 's to get the total $v_{\mathrm{a}, \mathrm{b}}$ :

$$
v_{\mathrm{a}, \mathrm{~b}}=v_{\mathrm{a}, \mathrm{~b} 1}+v_{\mathrm{a}, \mathrm{~b} 2}=v_{\mathrm{s}} \frac{R_{2}(1+\alpha)}{R_{1}+R_{2}(1+\alpha)}+i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

b) One consistency check is to set $v_{\mathrm{s}}=0 \mathrm{~V}$ and set $\alpha=1$, causing $R_{1}$ and the dependent source to be in parallel and have the same current, implying that they are the same resistance. In parallel, they have resistance $R_{1} / 2$.

In this case, our circuit will have the following output voltage:

$$
v_{\mathrm{a}, \mathrm{~b}}=i_{\mathrm{s}} \cdot \frac{R_{1}}{2} \| R_{2}=i_{\mathrm{s}} \cdot \frac{\frac{R_{1}}{2} R_{2}}{\frac{R_{1}}{2}+R_{2}}=i_{\mathrm{s}} \cdot \frac{R_{1} R_{2}}{R_{1}+2 R_{2}}
$$

Now we check what value our formula from part (a) gives:

$$
v_{\mathrm{a}, \mathrm{~b}}=0 \cdot \frac{R_{2}(1+1)}{R_{1}+R_{2}(1+1)}+i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+1) R_{2}+R_{1}}=i_{\mathrm{s}} \frac{R_{1} R_{2}}{2 R_{2}+R_{1}}
$$

This agrees with what we expect, so the consistency check is satisfied. Many other checks are possible.
c) The voltage $v_{\mathrm{a}, \mathrm{b}}$ found in part (a) is the Thevenin equivalent voltage, so all we need now is $R_{\mathrm{Th}}$. Perhaps the simplest way to find $R_{\mathrm{Th}}$ is to turn off the independent sources and connect a current source to the output. We then determine $v_{\mathrm{a}, \mathrm{b}}$ across the current source and use Ohm's law to find $R_{\mathrm{Th}}$. For the source, we could use a value of $i_{\mathrm{s}}$, in which case we have exactly Case II of the superposition from part (a). Our voltage will then be $v_{\mathrm{a}, \mathrm{b} 2}$. Thus, we have the following value for $R_{\mathrm{Th}}$ :

$$
R_{\mathrm{Th}}=\frac{v_{\mathrm{a}, \mathrm{~b} 2}}{i_{\mathrm{s}}}=\frac{i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}}{i_{\mathrm{s}}}=\frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

An alternative approach to finding $R_{\mathrm{Th}}$ is to use the short-circuit current, $i_{\mathrm{sc}}$, that flows from $\mathbf{a}$ to $\mathbf{b}$ when a wire is connected across those terminals. In that case, the voltage on the top and bottom rails is zero. This means there is no voltage drop across $R_{2}$, and we may ignore $R_{2}$. Also, we have voltage $-v_{\mathrm{s}}$ on the top end of $R_{1}$, giving the current for $i_{\mathrm{x}}$ directly:

$$
i_{\mathrm{x}}=-\frac{v_{\mathrm{s}}}{R_{1}}
$$

Now we can write a current summation for the top rail:

$$
-\frac{v_{\mathrm{s}}}{R_{1}}+\alpha\left(-\frac{v_{\mathrm{s}}}{R_{1}}\right)-i_{\mathrm{s}}+i_{\mathrm{sc}}=0 \mathrm{~A}
$$

or

$$
i_{\mathrm{sc}}=\frac{v_{\mathrm{s}}}{R_{1}}-\alpha\left(-\frac{v_{\mathrm{s}}}{R_{1}}\right)+i_{\mathrm{s}}=v_{\mathrm{s}} \frac{1+\alpha}{R_{1}}+i_{\mathrm{s}}
$$

Using this current, we find $R_{\mathrm{Th}}$ :

$$
R_{\mathrm{Th}}=\frac{v_{\mathrm{a}, \mathrm{~b}}}{i_{\mathrm{sc}}}=\frac{v_{\mathrm{s}} \frac{R_{2}(1+\alpha)}{R_{1}+R_{2}(1+\alpha)}+i_{\mathrm{s}} \frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}}{v_{\mathrm{s}} \frac{1+\alpha}{R_{1}}+i_{\mathrm{s}}}
$$

or

$$
R_{\mathrm{Th}}=\frac{v_{\mathrm{a}, \mathrm{~b}}}{i_{\mathrm{sc}}}=\frac{R_{1} R_{2}}{R_{1}+R_{2}(1+\alpha)}
$$

d) Maximum power is obtained by setting $R_{\mathrm{L}}=R_{\mathrm{Th}}$ :

$$
R_{\mathrm{L}}=R_{\mathrm{Th}}=\frac{R_{1} R_{2}}{(1+\alpha) R_{2}+R_{1}}
$$

